

APPROXIMATION OF FUNCTIONS FROM THE CLASS $C_{\beta, \infty}^{\psi}$ BY POISSON INTEGRALS IN THE UNIFORM METRIC

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We obtain asymptotic equalities for upper bounds of deviations of the Poisson integrals on the class of continuous functions $C_{\beta, \infty}^{\psi}$ in the metric of the space C .

1. Statement of the Problem and Auxiliary Assertions

Let $f(\cdot)$ be a 2π -periodic Lebesgue-summable function ($f \in L_1$). The Poisson integral of the function f is introduced (see [1, p. 154] or [2, p. 161]) as the function $P(\rho; f; x)$ defined by the equality

$$P(\rho; f; x) = \frac{1}{\pi} \int_{-\pi}^{\pi} f(t+x) \left\{ \frac{1}{2} + \sum_{k=1}^{\infty} \rho^k \cos kt \right\} dt, \quad 0 \leq \rho < 1.$$

Setting $\rho = e^{-1/\delta}$, we represent the Poisson integral in the form

$$P_{\delta}(f; x) = \frac{1}{\pi} \int_{-\pi}^{\pi} f(t+x) \left\{ \frac{1}{2} + \sum_{k=1}^{\infty} e^{-k/\delta} \cos kt \right\} dt, \quad \delta > 0.$$

In the present paper, we consider the class $C_{\beta, \infty}^{\psi}$ introduced by Stepanets (see, e.g., [3–6]), which is defined as follows: Assume that a function f belongs to L_1 and its Fourier series has the form

$$\frac{a_0}{2} + \sum_{k=1}^{\infty} (a_k \cos kx + b_k \sin kx).$$

Let $\psi(k)$ be an arbitrary function of a natural argument and let β be a fixed real number. If the series

$$\sum_{k=1}^{\infty} \frac{1}{\psi(k)} \left(a_k \cos \left(kx + \frac{\pi\beta}{2} \right) + b_k \sin \left(kx + \frac{\pi\beta}{2} \right) \right)$$

is the Fourier series of a certain function $\varphi \in L_1$, then φ is called the (ψ, β) -derivative of the function f and is denoted by $f_{\beta}^{\psi}(\cdot)$. Let L_{β}^{ψ} denote the subset of all functions $f \in L_1$ that have (ψ, β) -derivatives. If f belongs

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to L_β^ψ and f_β^ψ belongs to \mathfrak{N} , $\mathfrak{N} \subseteq L_1$, then one says that f belongs to $L_\beta^\psi \mathfrak{N}$. The subsets of continuous functions from L_β^ψ and $L_\beta^\psi \mathfrak{N}$ are denoted by C_β^ψ and $C_\beta^\psi \mathfrak{N}$, respectively. Further, if \mathfrak{N} coincides with the unit ball of the space L_∞ , i.e.,

$$\mathfrak{N} = \left\{ f_\beta^\psi \in L_\infty : \operatorname{ess\,sup}_t |f_\beta^\psi(t)| \leq 1 \right\},$$

then the classes $C_\beta^\psi \mathfrak{N}$ are denoted by $C_{\beta,\infty}^\psi$.

In the present paper, we study the asymptotic behavior of the quantity

$$\mathcal{E} \left(C_{\beta,\infty}^\psi; P_\delta \right)_C = \sup_{f \in C_{\beta,\infty}^\psi} \|f(\cdot) - P_\delta(f; \cdot)\|_C \tag{1}$$

as $\delta \rightarrow \infty$.

Following Stepanets [6, p. 198], we call the problem of finding asymptotic relations for quantity (1) as $\delta \rightarrow \infty$ the *Kolmogorov–Nicol’skii problem* for Poisson integrals on the class $C_{\beta,\infty}^\psi$ in the uniform metric.

Let \mathfrak{M} denote the set of functions $\psi(\cdot)$ that satisfy the conditions

$$\mathfrak{M} = \left\{ \psi(t) : \psi(t) > 0, \psi(t_1) - 2\psi((t_1 + t_2)/2) + \psi(t_2) \geq 0 \quad \forall t_1, t_2 \in [1, \infty), \lim_{t \rightarrow \infty} \psi(t) = 0 \right\}.$$

Let \mathfrak{M}' denote the set of functions $\psi \in \mathfrak{M}$ for which

$$\int_1^\infty \frac{\psi(t)}{t} dt < \infty.$$

Using the characteristics

$$\eta(t) = \eta(\psi; t) = \psi^{-1} \frac{\psi(t)}{2}, \quad \mu(t) = \mu(\psi; t) = \frac{t}{\eta(t) - t}, \tag{2}$$

where ψ^{-1} is the function inverse to ψ , one customarily considers (see, e.g., [5, p. 93] or [6, p. 160]) the following subsets of the set \mathfrak{M} :

$$\mathfrak{M}_0 = \{ \psi \in \mathfrak{M} : 0 < \mu(\psi; t) \leq K \quad \forall t \geq 1 \},$$

$$\mathfrak{M}_C = \{ \psi \in \mathfrak{M} : 0 < K_1 < \mu(\psi; t) \leq K_2 \quad \forall t \geq 1 \},$$

$$\mathfrak{M}_\infty = \{ \psi \in \mathfrak{M} : 0 < K \leq \mu(\psi; t) < \infty \quad \forall t \geq 1 \}.$$

Here and in what follows, K and K_i denote constants, generally speaking, different in different relations and dependent on ψ .

Note that, for functions $\psi \in \mathfrak{M}'_0$ ($\mathfrak{M}'_0 = \mathfrak{M}_0 \cap \mathfrak{M}'$) slowly decreasing to zero, i.e., for functions ψ such that

$$\int_1^{\infty} \psi(t) dt = \infty,$$

the Kolmogorov–Nikol’skii problem was solved in [7]. The aim of the present paper is to find asymptotic equalities for upper bounds of deviations of Poisson integrals on the classes $C_{\beta, \infty}^{\psi}$ for $\beta \in R$ in the cases where $\psi \in \mathfrak{M}_C$ and $\psi \in \mathfrak{M}_{\infty}$, i.e., for functions $\psi(t)$ that decrease to zero as $t \rightarrow \infty$ faster than the function $1/t$, which determines the order of saturation of the linear approximation method generated by the operator P_{δ} .

If the Fourier transform

$$\hat{\tau}(t) = \hat{\tau}_{\delta}(t) = \frac{1}{\pi} \int_0^{\infty} \tau(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \tag{3}$$

of the function $\tau(\cdot)$ defined by the equalities

$$\tau(u) = \tau_{\delta}(u; \psi) = \begin{cases} (1 - e^{-u}) \frac{\psi(1)}{\psi(\delta)}, & 0 \leq u \leq \frac{1}{\delta}, \\ (1 - e^{-u}) \frac{\psi(\delta u)}{\psi(\delta)}, & u \geq \frac{1}{\delta} \end{cases} \tag{4}$$

is summable on the entire number axis, i.e., the integral $A(\tau)$

$$A(\tau) = \int_{-\infty}^{\infty} |\hat{\tau}_{\delta}(t)| dt \tag{5}$$

is convergent, then, for any $f \in C_{\beta, \infty}^{\psi}$, the following equality holds at every point $x \in R$:

$$f(x) - P_{\delta}(f; x) = \psi(\delta) \int_{-\infty}^{+\infty} f_{\beta}^{\psi}\left(x + \frac{t}{\delta}\right) \hat{\tau}_{\delta}(t) dt, \quad \delta > 0. \tag{6}$$

Note that, relation (6) can be obtained by repeating the arguments used in [6, p. 183]. Thus, to find asymptotic equalities for quantity (1) as $\delta \rightarrow \infty$ in the case where $\psi \in \mathfrak{M}_C$, $\psi \in \mathfrak{M}_{\infty}$, and $\beta \in R$, it is necessary to find conditions under which the Fourier transform $\hat{\tau}(t)$ is summable on the entire number axis.

2. Asymptotic Equalities for Upper Bounds of Deviations of Poisson Integrals from Functions of the Class $C_{\beta, \infty}^{\psi}$ in the Uniform Metric

The following statement is true:

Theorem 1. *Suppose that $\psi \in \mathfrak{M}_C$, the function $g(u) = u\psi(u)$ is convex downward on $[b, \infty)$, $b \geq 1$, and*

$$\int_1^\infty \psi(u)du < \infty. \tag{7}$$

Then the following asymptotic equality holds as $\delta \rightarrow \infty$:

$$\mathcal{E}\left(C_{\beta,\infty}^\psi; P_\delta\right)_C = \frac{1}{\delta} \sup_{f \in C_{\beta,\infty}^\psi} \|f_0^{(1)}(x)\|_C + O\left(\frac{1}{\delta^2} \int_1^\delta t\psi(t)dt + \frac{1}{\delta} \int_\delta^\infty \psi(t)dt\right), \tag{8}$$

where $f_0^{(1)}$ is the (ψ, β) -derivative of the function f for $\psi(t) = 1/t$ and $\beta = 0$.

Prior to the proof of Theorem 1, we consider the following lemma:

Lemma 1. *Suppose that all conditions of Theorem 1 are satisfied. Then a Fourier transform $\hat{\tau}(t)$ of the form (3) for the function $\tau(u)$ defined by (4) is summable on the entire number axis, i.e., integral (5) is convergent.*

Proof of Lemma 1. We set $\tau(u) = \varphi(u) + v(u)$, where

$$\varphi(u) = \begin{cases} u \frac{\psi(1)}{\psi(\delta)}, & 0 \leq u < \frac{1}{\delta}, \\ u \frac{\psi(\delta u)}{\psi(\delta)}, & u \geq \frac{1}{\delta}, \end{cases} \tag{9}$$

$$v(u) = \begin{cases} (1 - e^{-u} - u) \frac{\psi(1)}{\psi(\delta)}, & 0 \leq u \leq \frac{1}{\delta}, \\ (1 - e^{-u} - u) \frac{\psi(\delta u)}{\psi(\delta)}, & u \geq \frac{1}{\delta}, \end{cases} \tag{10}$$

and verify that the Fourier transforms

$$\hat{\varphi}(t) = \hat{\varphi}_\delta(t) = \frac{1}{\pi} \int_0^\infty \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du, \tag{11}$$

$$\hat{v}(t) = \hat{v}_\delta(t) = \frac{1}{\pi} \int_0^\infty v(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \tag{12}$$

of the functions φ and v , respectively, are summable on the entire number axis. Thus, it is necessary to show that the following integrals are convergent:

$$A(\varphi) = \int_{-\infty}^\infty |\hat{\varphi}_\delta(t)| dt, \tag{13}$$

$$A(v) = \int_{-\infty}^{\infty} |\hat{v}_{\delta}(t)| dt. \tag{14}$$

First, we prove the convergence of integral (13). According to Theorem 1 in [8], for the convergence of the integral $A(\varphi)$ it is necessary and sufficient that the following integrals be convergent:

$$\int_0^{1/2} u |d\varphi'(u)|, \quad \int_{1/2}^{\infty} |u - 1| |d\varphi'(u)|,$$

$$\left| \sin \frac{\beta\pi}{2} \right| \int_0^{\infty} \frac{|\varphi(u)|}{u} du, \quad \int_0^1 \frac{|\varphi(1-u) - \varphi(1+u)|}{u} du.$$

It follows from (9) that

$$\varphi''(u) = 0, \quad u \in \left[0, \frac{1}{\delta}\right),$$

and

$$\psi(\delta) |d\varphi'(u)| \leq (2\delta |\psi'(\delta u)| + u\delta^2 \psi''(\delta u)) du, \quad \psi \in \mathfrak{M}, \quad \text{for } u \geq \frac{1}{\delta}. \tag{15}$$

Since

$$\int_0^{1/2} u |d\varphi'(u)| = \int_{1/\delta}^{1/2} u |d\varphi'(u)| \leq \int_{1/\delta}^{\infty} u |d\varphi'(u)|$$

and

$$\int_{1/2}^{\infty} |u - 1| |d\varphi'(u)| \leq \int_{1/\delta}^{\infty} u |d\varphi'(u)|,$$

we obtain an estimate for the integral

$$\int_{1/\delta}^{\infty} u |d\varphi'(u)|$$

on each of the intervals $[1/\delta, b/\delta)$ and $[b/\delta, \infty)$ (for $\delta > 2b$). Taking (15) into account, we get

$$\int_{1/\delta}^{b/\delta} u |d\varphi'(u)| \leq \frac{2\delta}{\psi(\delta)} \int_{1/\delta}^{b/\delta} u |\psi'(\delta u)| du + \frac{\delta^2}{\psi(\delta)} \int_{1/\delta}^{b/\delta} u^2 \psi''(\delta u) du.$$

Integrating both integrals on the right-hand side of the last inequality by parts and taking into account that $\psi(\delta u) \leq \psi(1)$ for $u \in [1/\delta, b/\delta)$, we get

$$\int_{1/\delta}^{b/\delta} u |d\varphi'(u)| \leq \frac{K_1}{\delta\psi(\delta)}.$$

Further, we show that the following relations are true:

$$\lim_{u \rightarrow \infty} u\psi(u) = 0, \tag{16}$$

$$\lim_{u \rightarrow \infty} u^2\psi'(u) = 0. \tag{17}$$

Since the function $g(u) = u\psi(u)$ is convex downward for $u \geq b \geq 1$, the following cases are possible: either

$$\lim_{u \rightarrow \infty} g(u) = 0,$$

or

$$\lim_{u \rightarrow \infty} g(u) = K > 0,$$

or

$$\lim_{u \rightarrow \infty} g(u) = \infty.$$

Let

$$\lim_{u \rightarrow \infty} g(u) = K > 0.$$

Then there exists $0 < K_1 < K$ such that, for all $u \geq 1$, one has $g(u) > K_1$ and, hence,

$$\psi(u) > \frac{K_1}{u},$$

which contradicts the fact that, according to condition (7), the function $\psi(u)$ is summable on $[1, \infty)$.

Now assume that

$$\lim_{u \rightarrow \infty} g(u) = \infty,$$

i.e., for any $M > 0$, there exists $N > 0$ such that $g(u) > M$ for all $u > N$. Then

$$\int_1^x \psi(u)du = \int_1^N \psi(u)du + \int_N^x \frac{g(u)}{u}du > K_2 + \int_N^x \frac{M}{u}du = K_2 + M(\ln x - \ln N).$$

We again arrive at a contradiction with the condition of the summability of the function $\psi(u)$ on the interval $[1, \infty)$. It follows from the results presented above that relation (16) is true.

We now prove relation (17). The function $g'(u)$ is summable on $[1, \infty)$, whence

$$\lim_{u \rightarrow \infty} \int_{u/2}^u g'(x) dx = 0.$$

Since, the function $g(u)$ is convex downward for $u \geq b \geq 1$, we conclude that the function $(-g'(u))$ does not increase for $u \geq b$, and, hence,

$$-\int_{u/2}^u g'(x) dx > -\left(u - \frac{u}{2}\right)(\psi(u) + u\psi'(u)) = -\frac{1}{2}(u\psi(u) + u^2\psi'(u)).$$

This and relation (16) yield (17).

Taking into account that the function $g(u)$, $u \geq b \geq 1$, is convex downward and using relations (16) and (17), we obtain

$$\int_{b/\delta}^{\infty} u |d\varphi'(u)| = \int_{b/\delta}^{\infty} u d\varphi'(u) = \lim_{u \rightarrow \infty} u\varphi'(u) - \frac{b}{\delta}\varphi'\left(\frac{b}{\delta}\right) + \varphi\left(\frac{b}{\delta}\right) = \frac{K}{\delta\psi(\delta)}.$$

Thus,

$$\int_0^{1/2} u |d\varphi'(u)| = O\left(\frac{1}{\delta\psi(\delta)}\right) \quad \text{and} \quad \int_{1/2}^{\infty} |u - 1| |d\varphi'(u)| = O\left(\frac{1}{\delta\psi(\delta)}\right) \quad \text{as } \delta \rightarrow \infty. \tag{18}$$

Further, taking into account relation (9) and the inequality

$$\int_1^{\infty} \psi(u) du \leq K,$$

we get

$$\int_0^{\infty} \frac{|\varphi(u)|}{u} du = \int_0^{\infty} \frac{\varphi(u)}{u} du = \frac{\psi(1)}{\delta\psi(\delta)} + \frac{1}{\delta\psi(\delta)} \int_1^{\infty} \psi(u) du = O\left(\frac{1}{\delta\psi(\delta)}\right).$$

Finally, we estimate the integral

$$\int_0^1 |\varphi(1 - u) - \varphi(1 + u)| \frac{du}{u}.$$

For this purpose, we represent this integral as a sum of two integrals:

$$\int_0^1 \frac{|\varphi(1 - u) - \varphi(1 + u)|}{u} du = \int_0^{1-1/\delta} \frac{|\varphi(1 - u) - \varphi(1 + u)|}{u} du + \int_{1-1/\delta}^1 \frac{|\varphi(1 - u) - \varphi(1 + u)|}{u} du. \tag{19}$$

We estimate the first term on the right-hand side of (19) by adding and subtracting the quantity $(-2u)$ under the modulus sign in the integrand. As a result, we get

$$\int_0^{1-1/\delta} \frac{|\varphi(1-u) - \varphi(1+u)|}{u} du = \int_0^{1-1/\delta} \frac{|\varphi(1-u) - \varphi(1+u) - 2u|}{u} du + O(1). \tag{20}$$

It follows from (9) that, for $u \in [0, 1 - 1/\delta]$, we have

$$1 - u = 1 - \frac{\psi(\delta)}{\psi(\delta(1-u))} \varphi(1-u), \quad 1 + u = 1 - \frac{\psi(\delta)}{\psi(\delta(1+u))} \varphi(1+u).$$

Then

$$\begin{aligned} & \int_0^{1-1/\delta} \frac{|\varphi(1-u) - \varphi(1+u) - 2u|}{u} du \\ & \leq \int_0^{1-1/\delta} |\varphi(1-u)| \left| 1 - \frac{\psi(\delta)}{\psi(\delta(1-u))} \right| \frac{du}{u} + \int_0^{1-1/\delta} |\varphi(1+u)| \left| 1 - \frac{\psi(\delta)}{\psi(\delta(1+u))} \right| \frac{du}{u}. \end{aligned}$$

Since the function $\varphi(\cdot)$ satisfies the conditions of Lemma 2 in [8], we have

$$|\varphi(u)| \leq |\varphi(0)| + |\varphi(1)| + \int_0^{1/2} u |d\varphi'(u)| + \int_{1/2}^\infty |u-1| |d\varphi'(u)| := H(\varphi).$$

Thus,

$$\begin{aligned} & \int_0^{1-1/\delta} \frac{|\varphi(1-u) - \varphi(1+u) - 2u|}{u} du \\ & = H(\varphi) O \left(\int_0^{1-1/\delta} \frac{|\psi(\delta(1-u)) - \psi(\delta)|}{u\psi(\delta(1-u))} du + \int_0^{1-1/\delta} \frac{|\psi(\delta(1+u)) - \psi(\delta)|}{u\psi(\delta(1+u))} du \right). \end{aligned} \tag{21}$$

Taking into account relation (9) and estimates (18) and using (16), we get

$$H(\varphi) = O\left(\frac{1}{\delta\psi(\delta)}\right), \quad \delta \rightarrow \infty. \tag{22}$$

It was established in [7] that the following estimates hold for functions $\psi \in \mathfrak{M}_0$ as $\delta \rightarrow \infty$:

$$\int_0^{1-1/\delta} \frac{|\psi(\delta(1-u)) - \psi(\delta)|}{u\psi(\delta(1-u))} du = O(1), \quad \int_0^{1-1/\delta} \frac{|\psi(\delta(1+u)) - \psi(\delta)|}{u\psi(\delta(1+u))} du = O(1);$$

these estimates are also true for functions $\psi \in \mathfrak{M}_C$.

Combining relations (20)–(22), we get

$$\int_0^{1-1/\delta} \frac{|\varphi(1-u) - \varphi(1+u)|}{u} du = O\left(\frac{1}{\delta\psi(\delta)}\right).$$

By analogy, one can easily verify that the same estimate holds for the second term on the right-hand side of (19). Therefore,

$$\int_0^1 |\varphi(1-u) - \varphi(1+u)| \frac{du}{u} = O\left(\frac{1}{\delta\psi(\delta)}\right), \quad \delta \rightarrow \infty.$$

Thus, we have established the convergence of integral (13) in the case where $\psi \in \mathfrak{M}$, the function $g(u) = u\psi(u)$ is convex downward on $[b, \infty)$, $b \geq 1$, and condition (7) is satisfied. Let us prove the convergence of integral (14). To this end, by virtue of Theorem 1 in [8], it is necessary to estimate the integrals

$$\int_0^{1/2} u|dv'(u)|, \quad \int_{1/2}^{\infty} |u-1||dv'(u)|, \tag{23}$$

$$\left| \sin \frac{\beta\pi}{2} \right| \int_0^{\infty} \frac{|v(u)|}{u} du, \quad \int_0^1 \frac{|v(1-u) - v(1+u)|}{u} du, \tag{24}$$

where $v(u)$ is the function given by (10), which is defined and continuous for all $u \geq 0$.

To estimate the first integral in (23), we divide the segment $[0; 1/2]$ into the two parts $[0; 1/\delta]$ and $[1/\delta; 1/2]$. It follows from (10) that

$$v''(u) = -e^{-u} \frac{\psi(1)}{\psi(\delta)} \quad \text{for } u \in \left[0, \frac{1}{\delta}\right).$$

Therefore,

$$\int_0^{1/\delta} u|dv'(u)| = \frac{\psi(1)}{\psi(\delta)} \int_0^{1/\delta} ue^{-u} du \leq \frac{\psi(1)}{\psi(\delta)} \int_0^{1/\delta} u du = O\left(\frac{1}{\delta^2\psi(\delta)}\right). \tag{25}$$

It also follows from relation (10) and properties of the function $\psi \in \mathfrak{M}$ that, for $u \geq 1/\delta$, one has

$$|dv'(u)| \leq \left\{ |\bar{v}(u)| \frac{\delta^2\psi''(\delta u)}{\psi(\delta)} + 2|\bar{v}'(u)| \frac{\delta|\psi'(\delta u)|}{\psi(\delta)} + |\bar{v}''(u)| \frac{\psi(\delta u)}{\psi(\delta)} \right\} du, \tag{26}$$

where $\bar{v}(u) = 1 - e^{-u} - u$. Using the inequalities

$$|\bar{v}(u)| \leq \frac{u^2}{2}, \quad |\bar{v}'(u)| \leq u, \quad |\bar{v}''(u)| \leq 1,$$

we rewrite relation (26) in the form

$$|dv'(u)| \leq \left\{ u^2 \frac{\delta^2 \psi''(\delta u)}{2\psi(\delta)} + 2u \frac{\delta |\psi'(\delta u)|}{\psi(\delta)} + \frac{\psi(\delta u)}{\psi(\delta)} \right\} du, \quad \psi \in \mathfrak{M}. \tag{27}$$

Using (27), we obtain the following relation for the first integral in (23) on the segment $[1/\delta; 1/2]$:

$$\int_{1/\delta}^{1/2} u |dv(u)| \leq \frac{1}{\psi(\delta)} \int_{1/\delta}^{1/2} \frac{u^3}{2} \delta^2 \psi''(\delta u) du + \frac{2}{\psi(\delta)} \int_{1/\delta}^{1/2} u^2 \delta |\psi'(\delta u)| du + \frac{1}{\psi(\delta)} \int_{1/\delta}^{1/2} u \psi(\delta u) du.$$

Taking the first integral on the right-hand side of the last inequality, we obtain

$$\int_{1/\delta}^{1/2} u |dv(u)| \leq \frac{1}{\psi(\delta)} \left. \frac{u^3}{2} \delta \psi'(\delta u) \right|_{1/\delta}^{1/2} + \frac{7}{2\psi(\delta)} \int_{1/\delta}^{1/2} u^2 \delta |\psi'(\delta u)| du + \frac{1}{\psi(\delta)} \int_{1/\delta}^{1/2} u \psi(\delta u) du. \tag{28}$$

Further, we use the following statements:

Proposition 1 [6, p. 161]. *A function $\psi \in \mathfrak{M}$ belongs to \mathfrak{M}_C if and only if the quantity*

$$\alpha(t) = \frac{\psi(t)}{t |\psi'(t)|}, \quad \psi'(t) = \psi'(t + 0),$$

satisfies the condition $0 < K_1 \leq \alpha(t) \leq K_2 \quad \forall t \geq 1$.

Proposition 2 [6, p. 175]. *A function $\psi \in \mathfrak{M}$ belongs to \mathfrak{M}_0 if and only if, for an arbitrary fixed number $c > 1$, there exists a constant K such that the following inequality holds for all $t \geq 1$:*

$$\frac{\psi(t)}{\psi(ct)} \leq K.$$

Using the conditions of Proposition 1, for $\psi \in \mathfrak{M}_C$ we get

$$\frac{1}{\psi(\delta)} \int_{1/\delta}^{1/2} u^2 \delta |\psi'(\delta u)| du \leq \frac{K}{\psi(\delta)} \int_{1/\delta}^{1/2} u \psi(\delta u) du.$$

Then, using (28) and taking into account Proposition 2 (which is also true for functions $\psi \in \mathfrak{M}_C$) and the inequality

$$\int_1^\delta u \psi(u) du \geq K,$$

we obtain

$$\int_{1/\delta}^{1/2} u |dv(u)| \leq K_1 + \frac{K_2}{\delta^2 \psi(\delta)} + \frac{K_3}{\psi(\delta)} \int_{1/\delta}^{1/2} u \psi(\delta u) du \leq \frac{K}{\delta^2 \psi(\delta)} \int_1^{\delta} u \psi(u) du. \tag{29}$$

Combining (25) and (29), we get

$$\int_0^{1/2} u |dv(u)| = O\left(\frac{1}{\delta^2 \psi(\delta)} \int_1^{\delta} u \psi(u) du\right), \quad \delta \rightarrow \infty. \tag{30}$$

Let us estimate the second integral in (23). For the function $\bar{v}(u) = 1 - e^{-u} - u$, we have $|\bar{v}(u)| \leq u$, $|\bar{v}'(u)| \leq 1$, and $|\bar{v}''(u)| = e^{-u}$. Taking this into account and using (26), we obtain the following relation for $\delta \geq 2$:

$$\begin{aligned} \int_{1/2}^{\infty} |u - 1| |dv'(u)| &\leq \int_{1/2}^{\infty} u |dv'(u)| \\ &\leq \frac{1}{\psi(\delta)} \int_{1/2}^{\infty} u e^{-u} \psi(\delta u) du + \frac{2\delta}{\psi(\delta)} \int_{1/2}^{\infty} u |\psi'(\delta u)| du + \frac{\delta^2}{\psi(\delta)} \int_{1/2}^{\infty} u^2 \psi''(\delta u) du. \end{aligned} \tag{31}$$

Let us estimate the first integral on the right-hand side of (31). Since the function $\psi(\delta u)$, $\delta \geq 2$, decreases for $u \in [1/2, \infty]$, taking Proposition 2 into account we get

$$\frac{1}{\psi(\delta)} \int_{1/2}^{\infty} u e^{-u} \psi(\delta u) du \leq \frac{\psi(\delta/2)}{\psi(\delta)} \int_{1/2}^{\infty} u e^{-u} du = O(1). \tag{32}$$

Integrating the third integral on the right-hand side of inequality (31) by parts and using equality (17) and Propositions 1 and 2, we obtain the following relation for the functions $\psi(\delta u) \in \mathfrak{M}_C$, $u \geq 1/2$, $\delta \geq 2$:

$$\begin{aligned} \frac{\delta^2}{\psi(\delta)} \int_{1/2}^{\infty} u^2 \psi''(\delta u) du &= \frac{\delta}{\psi(\delta)} \int_{1/2}^{\infty} u^2 d\psi'(\delta u) \\ &= \frac{\delta}{\psi(\delta)} \lim_{u \rightarrow \infty} u^2 \psi'(\delta u) + \frac{(\delta/2)|\psi'(\delta/2)|}{2\psi(\delta)} + \frac{2\delta}{\psi(\delta)} \int_{1/2}^{\infty} u |\psi'(\delta u)| du \\ &\leq K_1 + \frac{2\delta}{\psi(\delta)} \int_{1/2}^{\infty} u |\psi'(\delta u)| du. \end{aligned} \tag{33}$$

It follows from (31)–(33) that

$$\int_{1/2}^{\infty} |u - 1| |dv'(u)| \leq K_2 + \frac{4\delta}{\psi(\delta)} \int_{1/2}^{\infty} u |\psi'(\delta u)| du.$$

Integrating the integral on the right-hand side of the last relation again by parts and using relation (16) and Proposition 2, we obtain

$$\begin{aligned} \int_{1/2}^{\infty} |u - 1| |dv'(u)| &\leq K_3 + \frac{4}{\psi(\delta)} \int_{1/2}^{\infty} \psi(\delta u) du \\ &\leq K_3 + \frac{2\psi(\delta/2)}{\psi(\delta)} + \frac{4}{\psi(\delta)} \int_1^{\infty} \psi(\delta u) du \leq K_4 + \frac{4}{\delta\psi(\delta)} \int_{\delta}^{\infty} \psi(u) du. \end{aligned}$$

Thus, the following estimate holds as $\delta \rightarrow \infty$:

$$\int_{1/2}^{\infty} |u - 1| |dv'(u)| = O\left(1 + \frac{1}{\delta\psi(\delta)} \int_{\delta}^{\infty} \psi(u) du\right). \tag{34}$$

To estimate the first integral in (24), we divide the interval $[0; \infty)$ into the following three parts: $[0; 1/\delta]$, $[1/\delta; 1]$, and $[1; \infty)$. Taking into account the inequality

$$e^{-u} \leq 1 - u + \frac{u^2}{2}, \quad u \geq 0, \tag{35}$$

we obtain

$$\int_0^{1/\delta} \frac{|v(u)|}{u} du = \frac{\psi(1)}{\psi(\delta)} \int_0^{1/\delta} (-1 + e^{-u} + u) \frac{du}{u} \leq \frac{\psi(1)}{2\psi(\delta)} \int_0^{1/\delta} u du = O\left(\frac{1}{\delta^2\psi(\delta)}\right),$$

$$\int_{1/\delta}^1 \frac{|v(u)|}{u} du \leq \int_{1/\delta}^1 u \frac{\psi(\delta u)}{2\psi(\delta)} du = O\left(\frac{1}{\delta^2\psi(\delta)} \int_1^{\delta} u \psi(u) du\right),$$

$$\int_1^{\infty} \frac{|v(u)|}{u} du = \frac{1}{\psi(\delta)} \int_1^{\infty} \psi(\delta u) \left(\frac{e^{-u} - 1}{u} + 1\right) du \leq \frac{1}{\psi(\delta)} \int_1^{\infty} \psi(\delta u) du.$$

Hence,

$$\int_0^{\infty} \frac{|v(u)|}{u} du = O \left(\frac{1}{\delta^2 \psi(\delta)} \int_1^{\delta} u \psi(u) du + \frac{1}{\delta \psi(\delta)} \int_{\delta}^{\infty} \psi(u) du \right). \tag{36}$$

Let us estimate the second integral in (24). By analogy with the proof of relation (58) in [7], we obtain

$$\int_0^1 |v(1-u) - v(1+u)| \frac{du}{u} = \int_0^1 |\lambda(1-u) - \lambda(1+u)| \frac{du}{u} + O(H(v)), \tag{37}$$

where $\lambda(u) = e^{-u} + u$ and

$$H(v) = |v(0)| + |v(1)| + \int_0^{1/2} u |dv'(u)| + \int_{1/2}^{\infty} |u-1| |dv'(u)|.$$

Taking into account relations (10), (30), and (34) and the inequality

$$\frac{1}{\delta^2 \psi(\delta)} \int_1^{\delta} u \psi(u) du \geq \frac{1}{\delta^2 \psi(\delta)} \delta \psi(\delta) \int_1^{\delta} du \geq K,$$

we get

$$H(v) = O \left(\frac{1}{\delta^2 \psi(\delta)} \int_1^{\delta} u \psi(u) du + \frac{1}{\delta \psi(\delta)} \int_{\delta}^{\infty} \psi(u) du \right), \quad \delta \rightarrow \infty. \tag{38}$$

Since

$$\int_0^1 |\lambda(1-u) - \lambda(1+u)| \frac{du}{u} \leq K, \tag{39}$$

relations (37)–(39) yield the following estimate as $\delta \rightarrow \infty$:

$$\int_0^1 |v(1-u) - v(1+u)| \frac{du}{u} = O \left(\frac{1}{\delta^2 \psi(\delta)} \int_1^{\delta} u \psi(u) du + \frac{1}{\delta \psi(\delta)} \int_{\delta}^{\infty} \psi(u) du \right). \tag{40}$$

Thus, according to Theorem 1 in [8], integral (14) is also convergent.

Lemma 1 is proved.

Proof of Theorem 1. Lemma 1 states that, under the conditions of Theorem 1, the Fourier transform $\hat{\tau}(t)$ (3) of the function $\tau(u) = \varphi(u) + \nu(u)$ is summable on the entire number axis. Then, for any function $f \in C_{\beta, \infty}^{\psi}$, equality (6) holds at every point $x \in R$.

Using the integral representation (6), we represent quantity (1) in the form

$$\begin{aligned} \mathcal{E} \left(C_{\beta, \infty}^{\psi}; P_{\delta} \right)_C &= \sup_{f \in C_{\beta, \infty}^{\psi}} \left\| \psi(\delta) \int_{-\infty}^{+\infty} f_{\beta}^{\psi} \left(x + \frac{t}{\delta} \right) \hat{\tau}(t) dt \right\|_C \\ &= \sup_{f \in C_{\beta, \infty}^{\psi}} \left\| \psi(\delta) \int_{-\infty}^{+\infty} f_{\beta}^{\psi} \left(x + \frac{t}{\delta} \right) (\hat{\varphi}(t) + \hat{\nu}(t)) dt \right\|_C. \end{aligned}$$

Using (14), we obtain

$$\mathcal{E} \left(C_{\beta, \infty}^{\psi}; P_{\delta} \right)_C = \sup_{f \in C_{\beta, \infty}^{\psi}} \left\| \psi(\delta) \int_{-\infty}^{+\infty} f_{\beta}^{\psi} \left(x + \frac{t}{\delta} \right) \hat{\varphi}(t) dt \right\|_C + O(\psi(\delta)A(\nu)). \tag{41}$$

Repeating the arguments of [3], one can easily verify that the Fourier series of the function

$$f_{\varphi}(x) = \int_{-\infty}^{+\infty} f_{\beta}^{\psi} \left(x + \frac{t}{\delta} \right) \hat{\varphi}(t) dt$$

has the form

$$S[f_{\varphi}] = \sum_{k=1}^{\infty} \frac{k}{\delta} \frac{1}{\psi(\delta)} (a_k \cos kx + b_k \sin kx),$$

where a_k and b_k are the Fourier coefficients of the function f . Therefore,

$$\int_{-\infty}^{+\infty} f_{\beta}^{\psi} \left(x + \frac{t}{\delta} \right) \hat{\varphi}(t) dt = \frac{1}{\delta \psi(\delta)} f_0^{(1)}(x), \tag{42}$$

where $f_0^{(1)}(\cdot)$ is the (ψ, β) -derivative of the function $f(\cdot)$ in the Stepanets sense for $\psi(t) = 1/t$ and $\beta = 0$. Combining (41) and (42), we obtain

$$\mathcal{E} \left(C_{\beta, \infty}^{\psi}; P_{\delta} \right)_C = \frac{1}{\delta} \sup_{f \in C_{\beta, \infty}^{\psi}} \left\| f_0^{(1)}(x) \right\|_C + O(\psi(\delta)A(\nu)), \quad \delta \rightarrow \infty. \tag{43}$$

Using inequalities (2.14) and (2.15) from [8] and relations (30), (34), (36), (38), and (40), we obtain the following estimate for the integral $A(v)$:

$$A(v) = O\left(\frac{1}{\delta^2 \psi(\delta)} \int_1^{\delta} u \psi(u) du + \frac{1}{\delta \psi(\delta)} \int_{\delta}^{\infty} \psi(u) du\right), \quad \delta \rightarrow \infty.$$

This and relation (43) yield (8).

Theorem 1 is proved.

Examples of functions satisfying the conditions of Theorem 1 are functions $\psi \in \mathfrak{M}$ that have the following form for $t \geq 1$:

$$\psi(t) = \frac{1}{t} \ln^{\alpha}(t + K), \quad K > 0, \quad \alpha < -1; \quad \psi(t) = \frac{1}{t^r} \ln^{\alpha}(t + K);$$

$$\psi(t) = \frac{1}{t^r} \arctan t; \quad \psi(t) = \frac{1}{t^r} (K + e^{-t}), \quad r > 1, \quad K > 0, \quad \alpha \in R.$$

In the second part of the present paper, we find a solution of the Kolmogorov–Nicol’skii problem for Poisson integrals on the classes $C_{\beta, \infty}^{\psi}$ of continuous periodic functions in the case where ψ belongs to \mathfrak{M}_{∞} .

Theorem 2. *If ψ belongs to \mathfrak{M} , a function $g(u)$ is convex downward for $u \in [b, \infty)$, $b \geq 1$, and*

$$\int_1^{\infty} u^2 \psi(u) du < \infty, \tag{44}$$

then the following asymptotic equality holds as $\delta \rightarrow \infty$:

$$\mathcal{E}\left(C_{\beta, \infty}^{\psi}; P_{\delta}\right)_C = \frac{1}{\delta} \sup_{f \in C_{\beta, \infty}^{\psi}} \|f_0^{(1)}(x)\|_C + O\left(\frac{1}{\delta^2}\right), \tag{45}$$

where $f_0^{(1)}$ is the (ψ, β) -derivative of the function f for $\psi(t) = 1/t$ and $\beta = 0$.

The proof of Theorem 2 is based on the following auxiliary statement:

Lemma 2. *Suppose that all conditions of Theorem 2 are satisfied. Then an integral $A(\tau)$ of the form (5) is convergent.*

Proof of Lemma 2. To establish the convergence of the integral $A(\tau)$ we represent the function $\tau(\cdot)$ (4) as the sum of the functions $\varphi(\cdot)$ and $\nu(\cdot)$ defined by (9) and (10), respectively. We investigate the convergence of integral (13). To this end, we divide the set $(-\infty, \infty)$ into the two subsets $(-\infty, \delta) \cup (\delta, +\infty)$ and $[-\delta, \delta]$.

Let us estimate the integral $A(\varphi)$ for $|t| > \delta$. To this end, we consider the integral

$$\int_0^{\infty} \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du$$

on each of the intervals $[0; 1/\delta)$ and $[1/\delta; \infty)$:

$$\int_0^\infty \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du = \left(\int_0^{1/\delta} + \int_{1/\delta}^\infty\right) \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du. \tag{46}$$

It follows from (9) that

$$\varphi(0) = 0, \quad \varphi\left(\frac{1}{\delta}\right) = \frac{\psi(1)}{\delta\psi(\delta)}, \quad \text{and} \quad \varphi'(0) = \varphi'\left(\frac{1}{\delta} - 0\right) = \frac{\psi(1)}{\psi(\delta)} \quad \text{for } u \in \left[0, \frac{1}{\delta}\right).$$

Integrating the first integral on the right-hand side of equality (46) twice by parts and taking into account that $\varphi''(u) = 0$, $u \in [0, 1/\delta)$, we obtain

$$\int_0^{1/\delta} \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du = \frac{\psi(1)}{t\delta\psi(\delta)} \sin\left(\frac{t}{\delta} + \frac{\beta\pi}{2}\right) + \frac{\psi(1)}{t^2\psi(\delta)} \left(\cos\left(\frac{t}{\delta} + \frac{\beta\pi}{2}\right) - \cos\frac{\beta\pi}{2}\right). \tag{47}$$

By virtue of the convexity of the function $g(u)$ and condition (44), relations (16) and (17) are true. For $u \geq 1/\delta$, we get

$$\begin{aligned} &\int_{1/\delta}^\infty \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \\ &= -\frac{\psi(1)}{t\delta\psi(\delta)} \sin\left(\frac{t}{\delta} + \frac{\beta\pi}{2}\right) - \frac{1}{t^2} \left(\frac{\psi(1)}{\psi(\delta)} + \frac{\psi'(1)}{\psi(\delta)}\right) \cos\left(\frac{t}{\delta} + \frac{\beta\pi}{2}\right) \\ &\quad - \frac{1}{t^2} \int_{1/\delta}^\infty \varphi''(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du. \end{aligned} \tag{48}$$

Combining relations (46)–(48), we obtain

$$\begin{aligned} &\int_0^\infty \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \\ &= -\frac{1}{t^2\psi(\delta)} \left(\psi(1) \cos\frac{\beta\pi}{2} + \psi'(1) \cos\left(\frac{t}{\delta} + \frac{\beta\pi}{2}\right)\right) - \frac{1}{t^2} \int_{1/\delta}^\infty \varphi''(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du. \end{aligned}$$

Thus,

$$\left| \int_0^\infty \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| \leq \frac{K}{t^2\psi(\delta)} + \frac{1}{t^2} \int_{1/\delta}^\infty |\varphi''(u)| du. \tag{49}$$

Using relation (15) and taking into account that

$$\lim_{u \rightarrow \infty} \psi(u) = 0 \quad \text{and} \quad \lim_{u \rightarrow \infty} u\psi'(u) = 0,$$

we get

$$\frac{1}{t^2} \int_{1/\delta}^{\infty} |\varphi''(u)| du \leq -\frac{2}{t^2\psi(\delta)} \int_{1/\delta}^{\infty} d\psi(\delta u) + \frac{\delta}{t^2\psi(\delta)} \int_{1/\delta}^{\infty} u d\psi'(\delta u) = \frac{3\psi(1) - \psi'(1)}{t^2\psi(\delta)}.$$

Using this relation and (49), we obtain

$$\left| \int_0^{\infty} \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| \leq \frac{K_1}{t^2\psi(\delta)},$$

whence

$$\int_{|t| \geq \delta} \left| \int_0^{\infty} \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt \leq \frac{2K_1}{\delta\psi(\delta)}. \tag{50}$$

Let us estimate the integral $A(\varphi)$ on the segment $[-\delta, \delta]$. Since condition (44) is satisfied, we have

$$\begin{aligned} & \int_{-\delta}^{\delta} \left| \int_0^{\infty} \varphi(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt \\ & \leq 2\delta \int_0^{\infty} |\varphi(u)| du = \frac{\psi(1)}{\delta\psi(\delta)} + \frac{2\delta}{\psi(\delta)} \int_{1/\delta}^{\infty} u\psi(\delta u) du \\ & = \frac{\psi(1)}{\delta\psi(\delta)} + \frac{2}{\delta\psi(\delta)} \int_1^{\infty} u\psi(u) du \leq \frac{K_2}{\delta\psi(\delta)}. \end{aligned} \tag{51}$$

Using relations (50) and (51), we conclude that the following estimate holds as $\delta \rightarrow \infty$:

$$A(\varphi) = O\left(\frac{1}{\delta\psi(\delta)}\right).$$

Thus, the transform $\hat{\varphi}(t)$ (11) is summable on the entire number axis.

We now establish the convergence of the integral $A(\nu)$ [see (14)], where $\hat{\nu}(t)$ is the transform (12) of the function $\nu(\cdot)$ defined by relation (10). To this end, we divide the set $(-\infty, \infty)$ into the two parts $[-\delta, \delta]$ and $|t| > \delta$ so that

$$A(v) = \frac{1}{\pi} \int_{-\delta}^{\delta} \left| \int_0^{\infty} v(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt + \frac{1}{\pi} \int_{|t|>\delta} \left| \int_0^{\infty} v(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt := I_1 + I_2. \tag{52}$$

Let us estimate the integral

$$I_1 = \frac{1}{\pi} \int_{-\delta}^{\delta} \left| \int_0^{\infty} v(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt.$$

We have

$$I_1 \leq \frac{1}{\pi} \int_{-\delta}^{\delta} \left| \int_0^{1/\delta} v(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt + \frac{1}{\pi} \int_{-\delta}^{\delta} \left| \int_{1/\delta}^{\infty} v(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt. \tag{53}$$

Taking inequality (35) into account, we obtain

$$\begin{aligned} \frac{1}{\pi} \int_{-\delta}^{\delta} \left| \int_0^{1/\delta} v(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt &\leq \frac{1}{\pi} \int_{-\delta}^{\delta} \int_0^{1/\delta} |v(u)| du dt \\ &= \frac{2\delta\psi(1)}{\pi\psi(\delta)} \int_0^{1/\delta} (e^{-u} + u - 1) du \leq \frac{\psi(1)}{3\pi\delta^2\psi(\delta)}. \end{aligned} \tag{54}$$

According to the conditions of Lemma 2, we have

$$\int_1^{\infty} u^2 \psi(u) du < \infty.$$

Using inequality (35) once again, we obtain the following estimate for the second integral on the right-hand side of (53):

$$\begin{aligned} \frac{1}{\pi} \int_{-\delta}^{\delta} \left| \int_{1/\delta}^{\infty} v(u) \cos\left(ut + \frac{\beta\pi}{2}\right) du \right| dt &\leq \frac{1}{\pi} \int_{-\delta}^{\delta} \int_{1/\delta}^{\infty} |v(u)| du dt = \frac{2\delta}{\pi\psi(\delta)} \int_{1/\delta}^{\infty} (e^{-u} + u - 1) \psi(\delta u) du \\ &\leq \frac{\delta}{\pi\psi(\delta)} \int_{1/\delta}^{\infty} u^2 \psi(\delta u) du \leq \frac{K}{\pi\delta^2\psi(\delta)}. \end{aligned} \tag{55}$$

It follows from relations (53)–(55) that

$$I_1 = O\left(\frac{1}{\delta^2\psi(\delta)}\right), \quad \delta \rightarrow \infty. \tag{56}$$

Let us estimate the integral

$$I_2 = \frac{1}{\pi} \int_{|t| > \delta} \left| \int_0^{\infty} v(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du \right| dt.$$

Integrating twice by parts and taking into account that $v(0) = 0$ and $v'(0) = 0$, we get

$$\begin{aligned} & \int_0^{1/\delta} v(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du \\ &= \frac{1}{t} v \left(\frac{1}{\delta} \right) \sin \left(\frac{t}{\delta} + \frac{\beta\pi}{2} \right) + \frac{1}{t^2} v' \left(\frac{1}{\delta} - 0 \right) \cos \left(\frac{t}{\delta} + \frac{\beta\pi}{2} \right) - \frac{1}{t^2} \int_0^{1/\delta} v''(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du. \end{aligned} \quad (57)$$

Taking into account that

$$\lim_{u \rightarrow \infty} v(u) = 0 \quad \text{and} \quad \lim_{u \rightarrow \infty} v'(u) = 0,$$

which follows from (16) and (17), we obtain

$$\begin{aligned} & \int_{1/\delta}^{\infty} v(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du \\ &= -\frac{1}{t} v \left(\frac{1}{\delta} \right) \sin \left(\frac{t}{\delta} + \frac{\beta\pi}{2} \right) - \frac{1}{t^2} v' \left(\frac{1}{\delta} \right) \cos \left(\frac{t}{\delta} + \frac{\beta\pi}{2} \right) - \frac{1}{t^2} \int_{1/\delta}^{\infty} v''(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du. \end{aligned} \quad (58)$$

Combining (57) and (58), we get

$$\begin{aligned} & \int_0^{\infty} v(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du \\ &= \frac{1}{t^2} \left(\frac{1}{\delta} + e^{-1/\delta} - 1 \right) \frac{\delta\psi'(1)}{\psi(\delta)} \cos \left(\frac{t}{\delta} + \frac{\beta\pi}{2} \right) - \frac{1}{t^2} \left[\int_0^{1/\delta} + \int_{1/\delta}^{\infty} \right] v''(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du. \end{aligned}$$

Using inequality (35), we obtain

$$\left| \int_0^{\infty} v(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du \right| \leq \frac{1}{t^2} \left(\frac{K}{\delta\psi(\delta)} + \int_0^{1/\delta} |v''(u)| du + \int_{1/\delta}^{\infty} |v''(u)| du \right). \quad (59)$$

Since

$$v''(u) = -e^{-u} \frac{\psi(1)}{\psi(\delta)} \quad \text{for } u \in \left[0, \frac{1}{\delta}\right],$$

we get

$$\frac{1}{t^2} \int_0^{1/\delta} |v''(u)| du = \frac{\psi(1)}{t^2 \psi(\delta)} \int_0^{1/\delta} e^{-u} du \leq \frac{\psi(1)}{t^2 \delta \psi(\delta)}. \tag{60}$$

For the estimation of the second integral on the right-hand side of (59), we use relations (27), (16), and (17). As a result, we obtain

$$\begin{aligned} \frac{1}{t^2} \int_{1/\delta}^{\infty} |v''(u)| du &\leq \frac{1}{t^2 \psi(\delta)} \left(\int_{1/\delta}^{\infty} \psi(\delta u) du - 2 \int_{1/\delta}^{\infty} u d\psi(\delta u) + \frac{\delta}{2} \int_{1/\delta}^{\infty} u^2 d\psi'(\delta u) \right) \\ &= \frac{1}{t^2 \psi(\delta)} \left(\int_{1/\delta}^{\infty} \psi(\delta u) du - 2 \left(\lim_{u \rightarrow \infty} u \psi(\delta u) - \frac{\psi(1)}{\delta} - \int_{1/\delta}^{\infty} \psi(\delta u) du \right) \right. \\ &\quad \left. + \frac{\delta}{2} \left(\lim_{u \rightarrow \infty} u^2 \psi'(\delta u) - \frac{\psi'(1)}{\delta^2} \right) - \int_{1/\delta}^{\infty} u d\psi(\delta u) \right) \\ &= \frac{1}{t^2 \psi(\delta)} \left(4 \int_{1/\delta}^{\infty} \psi(\delta u) du + \frac{3\psi(1)}{\delta} - \frac{\psi'(1)}{2\delta} \right). \end{aligned}$$

Since

$$\int_1^{\infty} \psi(u) du < \infty,$$

we have

$$\frac{1}{t^2} \int_{1/\delta}^{\infty} |v''(u)| du \leq \frac{K}{t^2 \delta \psi(\delta)}. \tag{61}$$

It follows from relations (59)–(61) that

$$\left| \int_0^{\infty} v(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du \right| \leq \frac{K_1}{t^2 \delta \psi(\delta)}.$$

Then the following relation holds as $\delta \rightarrow \infty$:

$$I_2 = \frac{1}{\pi} \int_{|t|>\delta} \left| \int_0^{\infty} v(u) \cos \left(ut + \frac{\beta\pi}{2} \right) du \right| dt = O \left(\frac{1}{\delta^2 \psi(\delta)} \right). \tag{62}$$

Combining relations (52), (56), and (62), we get

$$A(v) = O \left(\frac{1}{\delta^2 \psi(\delta)} \right), \quad \delta \rightarrow \infty. \tag{63}$$

Lemma 2 is proved.

Proof of Theorem 2. Lemma 2 states that integrals (13) and (14) are summable under the conditions of Theorem 2. Therefore, using relation (43) and taking estimate (63) into account, we obtain equality (45).

Theorem 2 is proved.

Examples of functions satisfying the conditions of Theorem 2 are functions $\psi \in \mathfrak{M}$ that have the following form for $t \geq 1$:

$$\psi(t) = \frac{\ln^{\alpha}(t + K)}{t^r}, \quad \psi(t) = \frac{1}{t^r}(K + e^{-t}), \quad r > 3, \quad K > 0, \quad \alpha \in R,$$

$$\psi(t) = t^r e^{-Kt^{\alpha}}, \quad \psi(t) = \ln^r(t + e)e^{-Kt^{\alpha}}, \quad K > 0, \quad \alpha > 0, \quad r \in R.$$

Assume that a function $\mu(\cdot)$ is associated with a function $\psi \in \mathfrak{M}$ by relation (2). Theorem 2 yields the following corollary:

Corollary 1. If ψ belongs to \mathfrak{M}_{∞} , the function $g(u)$ is convex downward for $u \in [b, \infty)$, $b \geq 1$, and

$$\lim_{t \rightarrow \infty} \mu(\psi; t) = \infty, \tag{64}$$

then the asymptotic equality (45) holds as $\delta \rightarrow \infty$.

Proof. It suffices to verify that condition (64) guarantees the convergence of the integral

$$\int_1^{\infty} u^2 \psi(u) du.$$

It follows from relations (12.24) in [6, p. 164] that the following inequality holds for any function $\psi \in \mathfrak{M}$:

$$\frac{\psi(t)}{|\psi'(t)|} \leq 2(\eta(t) - t) \quad \forall t \geq 1. \tag{65}$$

In view of (65), for any $r \geq 0$ one has

$$(t^r \psi(t))' = rt^{r-1} \psi(t) - t^r |\psi'(t)| \leq t^r |\psi'(t)| \left(2r \frac{\eta(t) - t}{t} - 1 \right). \quad (66)$$

According to (64), the value $(\eta(t) - t)/t$ tends to zero as $t \rightarrow \infty$. Using relations (66), we conclude that, for any $r \geq 0$, there exists a number $t_0 = t_0(r, \psi)$ such that the function $t^r \psi(t)$ does not increase for $t > t_0$. Then

$$\int_1^{\infty} u^2 \psi(u) du = \int_1^{\infty} \frac{u^4 \psi(u)}{u^2} du \leq K \int_1^{\infty} \frac{du}{u^2} < \infty.$$

Thus, all conditions of Theorem 2 are satisfied. Therefore, equality (45) is true.

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